

the scale factor is $\theta=1+2\cos 2\pi/7 \approx 2.247$ if we take $l_T/l_S=2\cos 2\pi/7$ and $l_T/l_S=1-2\cos 3\pi/7$. We can get rational approximants by defining the numbers

$$L_{n+1}=L_n+S_n+T_n, S_{n+1}=L_n+S_n, T_{n+1}=L_n$$

with $L_0=1, S_0=0, T_0=0$ or also $L_0=0, S_0=1, T_0=0$.

The temporal structure derived from the substitution rules

$$L \rightarrow LSL, S \rightarrow LST, T \rightarrow ST$$

is also selfsimilar with scaling factor $\phi=1+2\cos\pi/7$. Observe that the word selfsimilarity is not being used in the same sense as for fractals. In fact there is always a minimum separation between points; there is not the ever finer detail that occurs in fractals. Obviously, infinite sequences of regularly spaced impulses are also selfsimilar, in the sense that if we replace each impulse by two impulses separated half the distance, the result is again an infinite sequence of regularly spaced impulses. But one of the most interesting properties of rhythms constructed with the aperiodic sequences considered in this work, is that they are not predictable: given a word it is not always possible to know what is the letter after a given one, without reproducing the whole word. Other substitutional sequences have been already used to produce musical shapes (see for instance [4]).

2 Time structures with discrete spectrum

A distribution of impulses on the points t_k of the time axis, with k integer, can be represented by the function

$$\rho(t) = \sum_k \delta(t - t_k) \quad (2)$$

where the Dirac delta-function $\delta(x)$ has the properties: $\delta(x)=0$ unless $x=0, \delta(0)=\infty$. Its Fourier transform is

$$\lim_{N \rightarrow \infty} (1/N) \sum_k \exp(it_k) \quad (3)$$

where N is the number of impulses. The Fourier transform of a sequence of impulses distributed along the time axis, is called the spectrum of the sequence, and is a sum of discrete and continuous components. The discrete component indicates order, the continuous component disorder. In this section we consider the aperiodic systems defined in section 1. The Fibonacci system and the LST-system with scaling factor θ , have discrete spectrum and therefore "enough order".

A periodic distribution with period of length a can be represented by the function

$$M_a(t) = \sum_k \delta(t - ka) \quad (4)$$

If t denotes the time in seconds, then the Fourier transform of $M_a(t)$ is proportional to $M_{2\pi/a}(\omega)$ where ω denotes the frequency in Hertz. For instance if we take $a=\pi/33$ we get the harmonic series: C₂, C₃, G₃, C₄, E₄, G₄, Bb₄, C₅, D₅, E₅... all of them with the same amplitude, which is measured by the proportionality factor of the function $M_{2\pi/a}(\omega)$.

In the Fibonacci case the impulse distribution is

$$M_a(t) = \sum_n \delta(t - t_n) \quad (5)$$

where t_n is given by equation (1). The spectrum of the distribution is discrete, and can be computed with the help of the golden number τ and two integers p and q , through the following expression:

$$\omega_{pq} = (2\pi / (1 + 1/\tau^2)) [p + q/\tau] \quad (6)$$

and with a different amplitude for each component. If we take $X = 2\pi q - \omega_{pq} / \tau$, then the amplitude is proportional to $(2 \sin(X/2)) / X$ (see Figure 1).

A pitch defined by p and q is more intense if $\tau q - p$ is small or p/q close to τ , that is when p, q are successive Fibonacci integers F_n, F_{n-1} . Outside this sequence, the amplitudes decrease strongly and, above a certain amplitude threshold, the number of partials below a fixed frequency is always finite.

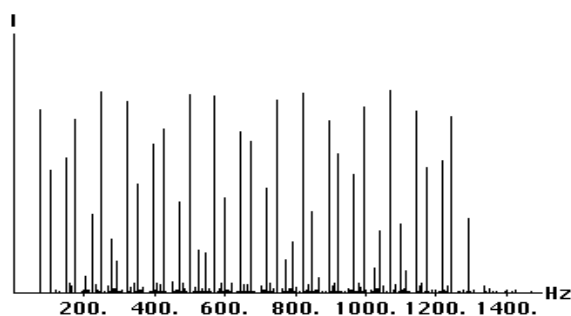


Figure 1. Spectrum for the Fibonacci system.

Given an arbitrary aperiodic system, it is not easy to find closed expressions for the amplitudes, as in the Fibonacci case. However recursion relations can be obtained allowing to get the Fourier amplitudes in an efficient way ([1] and references therein).

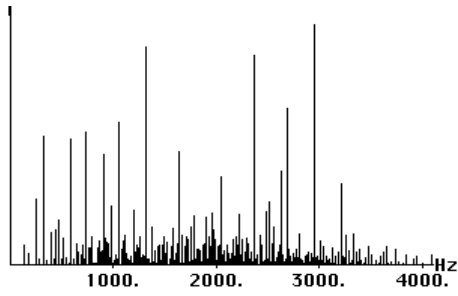


Figure 2. Discrete part of the spectrum for the L-S-T system with scaling factor θ

The order that lies behind the chain of temporal intervals, is reflected in the frequency distribution. In the AB-case it can be shown, that a set of partials separated by intervals of two lengths can be taken. The computations show that the ratio of the two intervals is the golden number (see the wavetable f1 in section.3, and compare the differences between consecutive partials). We can also extract from the spectrum of the first LST-system, a set of partials separated by intervals of three lengths, with the same ratio as l_L, l_S, l_T (see f2 in section.3).

The words obtained in the frequency space do not belong to the language generated by the grammar defining the temporal chains. In the Fibonacci words two B intervals can never be adjacent, nor can three A intervals. It can be seen that in the frequency space there are two consecutive short intervals.

3 Some examples

Due to the lack of translational symmetry, the spectra are always inharmonic. If the Csound language is used, the pairs frequency-amplitude for the Fibonacci system give the following composite waveform of weighted sums of simple sinusoids:

f1	0	4096	9	1	6.92784	0
2.618	19.4934	0	3.618	12.9569	0	5.236
14.398	0	6.236	18.5066	0	7.854	8.27371
0	8.854	21.363	0			
9.854	5.67165	0	11.472	20.3011	0	12.472
11.49	0	14.09	15.7818	0	15.09	17.3651
16.7	9.6877	0		17.71	21.022	0
4.27182	0	20.326	20.9093	0	21.326	
10.0286	0					

where we remind that the three p-fields in GEN09 correspond to the partial, strength of the partial and initial phase. A function table for the first LST system is:

f2	0	4096	9	1	0.659804	0
2.24734	0.875313	0	3.24751	1.08665	0	
4.49468	0.292446	0	5.04974	5.47652	0	

6.29692	1.37356	0	7.29708	1.94143	0
8.29725	0.46514	0	9.54442	0.28039	0
10.0995	0.918011	0	11.3467	5.61295	0
12.3468	1.19516	0	13.594	1.06184	0
14.149	4.70316	0	15.3962	2.48997	0
16.3964	6.05437	0	17.3966	1.28956	0

Bell-like timbres can be generated with the classical additive synthesis technique due to J.C.Risset [5]. The amplitude peaks of the partials are taken from f1 and f2 with exponential decay, durations inverse to their frequencies and beatings of the lowest two partials.

Other examples can be obtained from additive synthesis of filtered noises with center frequencies located at multiples of the partials given in f1 and f2, appropriate bandwidths, and gaussian envelopes with peaks placed at different times according with other parameters like frequency or amplitude.

If the recursion relations are used, temporal evolutions of the different partials can be implemented also, by increasing the number of impulses in the time axis: the ratio of amplitudes to number of time points stabilizes, when we increase the number of iterations.

4 Conclusion

In this work, aperiodic ordered temporal structures having discrete inharmonic spectra, have been discussed. A criterion based in number theory has been given in [6] to characterize a system with discrete spectrum. According with their results, the AB-system with scaling factor $\tau+2$ and the LST-system with scaling factor ϕ have no discrete part in their Fourier transforms.

In the compositional level, selfsimilarity suggest that musical form can also be articulated with the help of these multilevel hierarchies. The underlying grammar structure can represent connections between different sections of music having a high structural generality. Also the range of raw musical data that can be represented is high. The rhythmic structure shows both autonomy and solidarity with the pitch-intensity material.

Although only 1D examples have been considered in this work, the analysis of aperiodic structures in 2D and 3D [1] are also a rich source of musical applications. Some of them have discrete Fourier components and can be described in terms of formal grammars.

Spectral modeling of musical sounds (SMS) representing sinusoids and noise as two separate components has been developed in the last years (see [7] and references therein). The analysis detects time

